Superconducting La$_2$CuO$_{4+x}$ prepared by oxygenation at high pressure: A Raman-scattering study

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Superconducting crystals of La$_2$CuO$_{4+x}$ prepared by high-pressure oxygenation have been analyzed by Raman spectroscopy. A direct comparison of the role of excess oxygen was made by examining the same crystals with and without excess oxygen. La$_2$CuO$_{4+x}$, like nonsuperconducting La$_2$CuO$_{4}$, is found to have a soft phonon associated with a tetragonal-to-orthorhombic phase transition. The $A_g$ phonons of La$_2$CuO$_{4}$ and La$_2$CuO$_{4+x}$ occur at essentially the same frequency. At room temperature, La$_2$CuO$_{4+x}$ has no well-defined peak from two-magnon scattering, unlike La$_2$CuO$_{4}$. However, in its phase-separated form, La$_2$CuO$_{4+x}$ exhibits well-defined, two-magnon scattering. This establishes that the La$_2$CuO$_{4}$ phase present in La$_2$CuO$_{4+x}$ at low temperatures is antiferromagnetic. La$_2$CuO$_{4+x}$ samples prepared by slightly different methods are found to have differing amounts of excess oxygen, as indicated by variations in the intensity of the phonon and magnetic scattering. Certain samples of La$_2$CuO$_{4+x}$ had a phonon peak at 630 cm$^{-1}$ that is absent in La$_2$CuO$_{4}$. While a definitive assignment is not possible, the frequency of this peak is consistent with a peroxidelike species in La$_2$CuO$_{4+x}$.

I. INTRODUCTION

High-temperature superconductivity was first achieved by doping nonsuperconducting (NSC) La$_2$CuO$_4$ with Ba or Sr to give La$_{2-x}$R$_x$CuO$_4$ ($R =$ Ba or Sr).$^1$ Illustrated in Fig. 1(a) is the structure of the parent compound La$_2$CuO$_{4}$. It is possible, however, to convert La$_2$CuO$_{4}$ into a superconductor in another manner, namely, by putting excess oxygen into the structure.$^{2,3}$ When sufficient excess oxygen is introduced by heat treatment in high-pressure oxygen, a stable, bulk superconductor is obtained.$^4$ The resulting materials, La$_2$CuO$_{4+x}$, have average oxygen stoichiometries of about 4.03–4.04, superconducting transition temperatures, $T_c$, of up to 39 $K$, and superconducting volume fractions of over 50%.$^5$ From the viewpoint of charge balance, introducing excess anionic oxygen achieves the same effect as does substituting Ba or Sr (formal valence +2) for La (formal valence +3).

While it is accepted that Ba and Sr substitute for La on the La site,$^6,7$ the structural defect and the configuration formed by the excess oxygen of superconducting (SC) La$_2$CuO$_{4+x}$ is not well established. Similar uncertainty is associated with the valence and local bonding of the excess oxygen. Early measurements of oxygen content by gravimetric analysis and chemical titration suggested that the excess oxygen was in the form of superoxide, O$_2$$^-$$^-$$^-$. This view was supported by analysis using photoelectron spectroscopy.$^8$ Zhou, Sinha, and Goodenoough$^9$ argued that disparities between gravimetric and titration measurements could be explained by surface oxygen, a result that is consistent with more recent single-crystal studies.$^{10}$ Other research has suggested that the excess oxy-

 gen is in the form of peroxide, O$_2$$^-$$^-$$^-$. For example, neutron-diffraction measurements$^{9,10}$ of a single crystal of average composition La$_2$CuO$_{4.03}$ showed as one interpretation that the excess oxygen resulted in a short O–O bond of 1.59 Å, consistent with bond lengths of peroxides. Fig. 1(b) illustrates the structural position of the excess oxygen as suggested by the neutron-diffraction study. Like NSC La$_2$CuO$_{4.0}$ SC La$_2$CuO$_{4+x}$ has a tetragonal-to-orthorhombic (T–O) phase transition upon cooling. With further cooling, SC La$_2$CuO$_{4+x}$ undergoes phase separation into two distinct orthorhombic phases.$^{11}$ One of the two coexisting phases is known to contain the excess oxygen and to be superconducting. While the other phase is known to be essentially stoichiometric La$_2$CuO$_{4.0}$, little is known about its electrical or magnetic properties.

Here, Raman scattering has been used to examine single crystals of La$_2$CuO$_{4+x}$. A direct comparison was made between the same samples with excess oxygen (i.e., superconducting La$_2$CuO$_{4+x}$) and after removing the excess oxygen (i.e., nonsuperconducting La$_2$CuO$_{4.0}$). If the excess oxygen of SC La$_2$CuO$_{4+x}$ is indeed in the form of some dioxygen complex such as O$_2$$^-$$^-$$^-$ or O$_2$$^-$$^-$$^-$, then these species should have local vibrational modes associated with them. Since the vibrational frequency of dioxygen complexes is highly sensitive to the valence of the complex, then vibrational analysis of La$_2$CuO$_{4+x}$ should help clarify the nature of the excess oxygen. In addition, this study has investigated the phase transition and phase separation of SC La$_2$CuO$_{4+x}$ using Raman spectroscopy. A soft phonon mode is found to be associated with the T–O phase transition in SC La$_2$CuO$_{4+x}$. Finally, a comparison is made between light scattering from magnetic spin fluctuations (i.e., magnon scattering) from SC La$_2$CuO$_{4+x}$ and
II. EXPERIMENTAL

Superconducting crystals of La$_2$CuO$_{4+x}$ were prepared by heating nonsuperconducting single crystals of La$_2$CuO$_{4.0}$ (NSC LaCuO$_{4.0}$) in either 3 kbar of O$_2$ at 575°C for 12 h or in 3 kbar of O$_2$ at 550°C for 60 h. Crystals prepared by the former conditions, referred to as SC LaCuO$_{4+x}$ No. 1, exhibited the onset of superconductivity at 39 K. Crystals prepared in the latter manner, referred to as SC LaCuO$_{4+x}$ No. 2, exhibited the onset of superconductivity at 38 K. A direct comparison of the effect of the excess oxygen was made by examining the SC LaCuO$_{4+x}$ No. 2 crystals by Raman spectroscopy before and after removing the excess oxygen by heating at 700°C in one atmosphere of O$_2$. Raman measurements were made in a 180° backscattering geometry using either 488- or 514.5-nm excitation. All the results shown here are for 488-nm excitation. The majority of Raman measurements were made on the interior surfaces of the crystals exposed by mechanical fracture. All spectra of SC LaCuO$_{4+x}$ shown here are from fracture surfaces, while the NSC LaCuO$_{4.0}$ spectra are from surfaces prepared by mechanical polishing. Polarization geometries are referred to the axes of the tetragonal structure illustrated in Fig. 1. In (yy) polarization, the incident light propagates perpendicular to the c axis and the incident and scattered light are polarized perpendicular to the c axis. In (zz) polarization, the incident and scattered light are polarized parallel to the c axis.

III. RAMAN SCATTERING RESULTS OF SUPERCONDUCTING La$_2$CuO$_{4+x}$

Like La$_2$CuO$_{4.0}$, SC LaCuO$_{4+x}$ undergoes a tetragonal-to-orthorhombic (T-O) transition upon cooling (space group $I4/mmm$ to $Cmca$). The T-O transition temperature is strongly dependent on the exact oxygen content of the sample. For example, the T-O transition occurs at about 530 K for SC La$_2$CuO$_{4+x}$. In contrast, the SC LaCuO$_{4+x}$ samples of this study are tetragonal at room temperature (295 K) and the T-O transition occurs at about 280 K. The Raman-active phonons allowed in the $I4/mmm$ structure are two of $A_{1g}$ symmetry and two of $E_g$ symmetry (see Table I). In the upper Raman spectrum of Fig. 2, two $A_{1g}$ phonons are observed at 224 and 430 cm$^{-1}$ in tetragonal SC LaCuO$_{4+x}$ at 295 K. The orthorhombic $Cmca$ structure has additional Raman-active phonons (see Table I). The lower spectrum (77 K) of Fig. 2 shows the five $A_g$ phonons of orthorhombic SC LaCuO$_{4+x}$. The frequencies of the five $A_g$ phonons of orthorhombic SC LaCuO$_{4+x}$ are listed in Table II.

SC LaCuO$_{4+x}$ has an additional feature not present in NSC LaCuO$_{4.0}$; namely SC LaCuO$_{4+x}$ separates into two distinct orthorhombic phases upon cooling. The phase-separation temperature is strongly dependent upon the exact amount of excess oxygen. For example, the material used in the study of Jorgensen et al.[11] undergoes

FIG. 1. Structures of tetragonal (a) La$_2$CuO$_{4,0}$, and (b) La$_2$CuO$_{4.08}$ as determined by neutron diffraction. The excess oxygen occupies site O(4), giving a short O(3)—O(4) bond distance of 1.59 Å. Drawings adapted from Chaillout and co-workers (Refs. 5 and 10).
the $I4/mmm$ to $Cmca$ transition at about 430 K and phase separation into two orthorhombic phases occurs at about 320 K. In contrast, the samples of the present study contain a greater amount of excess oxygen and the phase separation occurs at a lower temperature, about 260 K.\textsuperscript{14} While originally one phase of the phase-separated material was thought to be $Fmmm$,\textsuperscript{11} both structures are now thought to be $Cmca$:\textsuperscript{5}

<table>
<thead>
<tr>
<th>Tetragonal</th>
<th>$I4/mmm$</th>
<th>$2A_g$</th>
<th>$2E_g$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Orthorhombic</td>
<td>$Cmca$*</td>
<td>$5A_g$</td>
<td>$4B_{1g}$</td>
</tr>
<tr>
<td>Orthorhombic</td>
<td>$Fmmm$*</td>
<td>$2A_g$</td>
<td>$2B_{3g}$</td>
</tr>
</tbody>
</table>

*Included for completeness since $Fmmm$ has been suggested as one of the two phases in phase-separated SC LaCuO$_{4+x}$ (Ref. 11).

Neutron-diffraction analysis\textsuperscript{5,10} of a crystal prepared in a similar manner as those of the present study has established that one of the two coexisting phases is essentially stoichiometric LaCuO$_{4.0}$. The other coexisting phase contains the excess oxygen and has approximate composition La$_2$CuO$_{4.05}$. Interestingly, the existence of these two coexisting phases is not apparent in the Raman data, even at liquid helium temperatures. Specifically, no splitting of the five $A_g$ phonons is observed in Fig. 2, as might be expected from two different phases with slightly different lattice parameters and vibrational force constants.

The orthorhombic $Cmca$ phases differ from the tetragonal $I4/mmm$ phase by essentially a nearly rigid rotation of the CuO$_6$ octahedra about the original [110] tetragonal axis.\textsuperscript{6} This rotation is a Raman-active vibration\textsuperscript{15} and exhibits soft-mode behavior in SC LaCuO$_{4+x}$ as in NSC La$_2$CuO$_{4.0}$ (see Sec. IV). As shown in Fig. 2, the $A_g$ phonon at 131 cm$^{-1}$ at 77 K has softened to 112 cm$^{-1}$ at 250 K. This phonon could only be traced to about 250 K; above this temperature the phonon was not observable. Figure 3 shows the temperature dependence of the soft phonons of SC LaCuO$_{4+x}$ and NC LaCuO$_{4.0}$. Since the soft $A_g$ mode is allowed only in orthorhombic SC LaCuO$_{4+x}$ but not in the tetragonal phase, the disappearance of the soft $A_g$ mode above about 250 K is consistent with the T-O transition observed by other techniques near 280 K.\textsuperscript{5,14}

IV. COMPARISON WITH NONSUPERCONDUCTING LaCuO$_{4.0}$

While NSC LaCuO$_{4.0}$ has been extensively studied by Raman spectroscopy, the literature is inconsistent. (See review articles by Ferraro and Maroni\textsuperscript{16} and Feile\textsuperscript{17} for a summary of the literature results.) The five allowed $A_g$ phonons in the 300-430 cm$^{-1}$ range are shown in Fig. 4, along with the results of other investigations. In general, the results of these investigations show softening of the $A_g$ phonons compared to those in SC LaCuO$_{4+x}$, with the $3B_{2g}$ mode being the most sensitive to the softening. This softening is consistent with the phase separation into two soft modes in the orthorhombic phase. The $131$ cm$^{-1}$ phonon at 77 K is the orthorhombic $Cmca$ soft mode, and the other orthorhombic soft mode is at about 250 K. The other $A_g$ phonons are not affected by the phase separation, and their frequencies are nearly the same as those in SC LaCuO$_{4+x}$.
modes of orthorhombic LaCuO$_{4.0}$ are given in Table II. The frequencies reported here agree extremely well (to within 3 cm$^{-1}$) with the results of Sugai.\textsuperscript{15} This agreement is convincing evidence that these are truly the five $A_g$ phonons of orthorhombic LaCuO$_{4.0}$. Figure 3 shows the temperature dependence of the soft $A_g$ phonon associated with the T-O transition. This soft phonon has been extensively studied by both Raman spectroscopy\textsuperscript{15,18} and inelastic neutron scattering.\textsuperscript{19}

Several comments comparing SC LaCuO$_{4+x}$ and NSC LaCuO$_{4.0}$ are in order. First, the temperature dependence of the soft $A_g$ mode is nearly identical for the two materials below about 250 K. However, the soft mode of SC LaCuO$_{4+x}$ essentially disappears above about 250 K, roughly the temperature at which phase separation and the T-O transition occurs. Second, the $A_g$ phonons of orthorhombic SC LaCuO$_{4+x}$ and orthorhombic NSC LaCuO$_{4.0}$ occur at essentially the same frequency (see Table II). This insensitivity to the excess oxygen is surprising in view of the rather large structural distortion [see Fig. 1(b)] that the majority [\textsuperscript{15}70\% (Ref. 5)] of the sample has undergone at low temperature. However, differences do exist in the relative intensities of the $A_g$ phonons. Figure 4 shows a comparison between the orthorhombic forms of NSC LaCuO$_{4.0}$ (upper spectrum), SC LaCuO$_{4+x}$ No. 1 (middle spectrum), and SC LaCuO$_{4+x}$ No. 2 (bottom spectrum). Relative to the axial-oxygen vibration [O(1) in Fig. 1(a)] at 427 cm$^{-1}$, the peak height of the soft phonon (~131 cm$^{-1}$ at 77 K) decreases from NSC LaCuO$_{4.0}$ to SC LaCuO$_{4+x}$ No. 1 to SC LaCuO$_{4+x}$ No. 2. The La vibration (228 cm$^{-1}$) displays the opposite trend, i.e., it increases in relative intensity. Since SC LaCuO$_{4+x}$ has a higher oxygen content than NSC LaCuO$_{4.0}$, it is clear that the relative intensity of the soft phonon decreases with increasing oxygen content. Further, as supported in the next section, it is believed that the soft phonon of SC LaCuO$_{4+x}$ No. 2 is weaker than that of SC LaCuO$_{4+x}$ No. 1 due to the higher overall oxygen content of the former material. It is interesting that the same trends observed here in relative phonon intensities as La$_2$Cu$_4$O$_{4.0}$ is doped with excess oxygen are also observed as La$_2$Cu$_4$O$_{4.0}$ is doped with strontium.\textsuperscript{15}

V. MAGNETIC SCATTERING

Raman spectroscopy has been extensively used to study spin fluctuations in the Cu-O based materials. In the parent compound La$_2$Cu$_4$O$_{8.0}$, two-magnon scattering is found as a broad but well-defined peak at about 3000 cm$^{-1}$ with $B_{1g}$ (tetragonal) symmetry. [See Fig. 5(a) and

![Figure 3](image3.png)

**FIG. 3.** Temperature dependence of the soft $A_g$ phonon of SC LaCuO$_{4+x}$ and NSC LaCuO$_{4.0}$. Open squares are from crystal SC LaCuO$_{4+x}$ No. 1 and solid squares are from crystal SC LaCuO$_{4+x}$ No. 2.

![Figure 4](image4.png)

**FIG. 4.** Comparison of NSC LaCuO$_{4.0}$ (upper spectrum), SC LaCuO$_{4+x}$ No. 1 (middle spectrum), and SC LaCuO$_{4+x}$ No. 2 (bottom spectrum) in polarization geometries permitting $A_g$ symmetries (orthorhombic $Cmca$). Spectra are offset for clarity.
Refs. 15 and 20]. Since the two-magnon peak is dominated
by short-range magnetic order, it is little affected by
temperature and persists above \( T_N \). As shown in Figs.
5(b) and 5(c) at 295 K neither SC \( \text{LaCuO}_{4+x} \) No. 1 nor SC
\( \text{LaCuO}_{4+x} \) No. 2 exhibit a well-defined two-magnon
peak. This is not surprising since Lyons and Fleury\(^{21}\)
have shown that the two-magnon scattering of
\( \text{YBa}_2\text{Cu}_3\text{O}_6+\) shifts to lower frequency and greatly
broadens upon doping the parent nonsuperconducting
compound \( \text{YBa}_2\text{Cu}_3\text{O}_6 \) with oxygen. For the supercon-
ductor \( \text{YBa}_2\text{Cu}_3\text{O}_{6.9} \) \( (T_c \approx 90 \text{ K}) \), there is no well-defined
peak from magnon scattering. Similar results are ob-
served in the TI-based system. That is, the antiferromag-
cetic insulator \( \text{TlYBa}_2\text{Cu}_3\text{O}_{7} \) has a well-defined two-
magnon peak but the doped superconductor
\( \text{TICaBa}_2\text{Cu}_3\text{O}_{7} \) does not.\(^{22}\)

SC \( \text{LaCuO}_{4+x} \) is more complicated since it separates
into two coexisting phases below about 260 K. Neutron-
diffraction analysis has shown that from a structural
standpoint, one of these two phases is essentially
\( \text{La}_2\text{CuO}_{4.6} \).\(^{5,11}\) The other phase contains the excess oxy-
gen and is known to be superconducting.\(^{11}\) Less is
known about the \( \text{La}_2\text{CuO}_{4.0} \) phase in phase-separated SC
\( \text{LaCuO}_{4+x} \), although it has been suggested to be nonsu-
perconducting.\(^{12}\) Figures 5(b) and 5(c) illustrate the mag-
netic scattering observed from SC \( \text{LaCuO}_{4+x} \) No. 1 and
SC \( \text{LaCuO}_{4+x} \) No. 2 at 77 K, well below the temperature
of phase separation. SC \( \text{LaCuO}_{4+x} \) No. 1 [Fig. 5(b)] has a
well-defined two-magnon peak at about 3000 cm\(^{-1}\). At
77 K, SC \( \text{LaCuO}_{4+x} \) No. 2 [Fig. 5(c)] has an extremely
weak two-magnon peak at about 3000 cm\(^{-1}\). Since the
phase with the excess oxygen is known to be supercon-
ducting, it will not exhibit a well-defined two-magnon
peak, as discussed above. Therefore, the two-magnon
peak observed in SC \( \text{LaCuO}_{4+x} \) at 77 K must result from
the \( \text{La}_2\text{CuO}_{4.0} \) phase and this phase must be antiferro-
magnetic and nonsuperconducting.

At 77 K, the two-magnon peak of SC \( \text{LaCuO}_{4+x} \) No. 1 is
intense, while the two-magnon peak of SC \( \text{LaCuO}_{4+x} \)
No. 2 is barely discernible. This implies that SC
\( \text{LaCuO}_{4+x} \) No. 2 has less of the antiferromagnetic phase
\( \text{La}_2\text{CuO}_{4.0} \) than does SC \( \text{LaCuO}_{4+x} \) No. 1. Since the
amount of the \( \text{La}_2\text{CuO}_{4.0} \) phase that forms is reduced as
the average oxygen content of the material increases,\(^{11}\) it
is concluded that the SC \( \text{LaCuO}_{4+x} \) No. 2 material has a
higher oxygen content than does SC \( \text{LaCuO}_{4+x} \) No. 1.
This conclusion is consistent with the observation of Sec.
IV, where the intensity of the soft \( A \) phonon was found
to be related to the amount of excess oxygen.
Specifically, in Fig. 4, the soft phonon of SC \( \text{LaCuO}_{4+x} \)
No. 2 (bottom spectrum) was weaker than that of SC
\( \text{LaCuO}_{4+x} \) No. 1 (middle spectrum).

As seen in Fig. 5, both SC \( \text{LaCuO}_{4+x} \) and NSC \( \text{LaCuO}_{4.0} \)
have numerous Raman features between about
700 and 1500 cm\(^{-1}\). For NSC \( \text{LaCuO}_{4.0} \) these high-
frequency peaks have been interpreted as either second-
order phonons enhanced by some resonance process\(^{15,23}\)
or magnetic scattering from a spin-density wave in the
antiferromagnetic structure.\(^{24}\) These peaks are largely
absent in tetragonal \( \text{La}_2\text{CuO}_{4.0} \) and can be largely el-
minated by heat treatment in low pressures of oxygen.\(^{25}\)
For SC \( \text{LaCuO}_{4+x} \), this high-frequency structure is larg-
ely absent at 295 K [Figs. 5(b), 5(c), and 6] but is present at
lower temperatures. The intensity of the structure be-
tween 700 and 1500 cm\(^{-1}\) varies greatly from sample to
sample. Compare, for example, Figs. 5(b) (SC \( \text{LaCuO}_{4+x} \)
No. 1) and 5(c) (SC \( \text{LaCuO}_{4+x} \) No. 2). Even within a
sample, the intensity of this structure was variable, for
example, the spectra of Figs. 6 and 5(c), taken from
different spots on the same crystal.

\[ \text{LaCuO}_{4+x} \]

**FIG. 5.** Comparison of two-magnon scattering in (a) NSC
\( \text{LaCuO}_{4.0} \), (b) SC \( \text{LaCuO}_{4+x} \) No. 1, and (c) SC \( \text{LaCuO}_{4+x} \) No. 2. Spectra were obtained from the same region of each crystal
as the corresponding spectra of Fig. 4. The 77-K spectra are
offset from the 295-K spectra for clarity. The short horizontal
lines in (b) and (c) represent zero intensity for the 77-K spectra.
The notation \( x(yy)x \) refers to polarization of both the incident
and scattered light along the \( b \) = \( y \) axis with the incident laser
beam and the scattered light propagating along the \( a \) = \( x \) axis.
\( x(yy)x \) polarization geometry permits \( A_{1g} + B_{1g} \) symmetry in
tetragonal \( I4/mmm \) and \( A_2 + B_{2g} \) symmetry in orthorhombic
\( Cmca \). \( x(yy)x \) polarization geometry permits \( E_g \) symmetry in
tetragonal \( I4/mmm \) and \( B_{1g} + B_{3g} \) symmetry in orthorhombic
\( Cmca \).
observed using \((yy)\) polarization, examination of the same spot using \((zz)\) polarization revealed that the soft \(A_g\) phonon was weaker than in any of the spectra of Fig. 4. This includes the bottom spectra of Fig. 4, obtained from a region of SC \(\text{LaCuO}_4^{+ \pm x}\) No. 2 where the 630-cm\(^{-1}\) peak was not observed. Given the relationship between the intensity of the soft phonon and oxygen content, we conclude that the 630-cm\(^{-1}\) peak is associated with regions having the highest amounts of excess oxygen. The absence of the 630-cm\(^{-1}\) peak within some regions of the SC \(\text{LaCuO}_4^{+ \pm x}\) No. 2 crystals indicates that the crystals are not uniformly loaded with oxygen. This is not surprising, given the size of the crystals (about 1–2 mm in thickness).

There are two sources of Raman-scattering peaks in SC \(\text{LaCuO}_4^{+ \pm x}\): phonon scattering and magnetic scattering. (An additional source of intensity is electronic scattering from charge carriers. However, this scattering, if present, should be in the form of very broad features, i.e., a continuum.) It is unlikely that the peak at 630 cm\(^{-1}\) results from magnetic scattering. Below its Neel temperature \(T_N\) of \(\approx 300\) K,\(^{15}\) SC \(\text{LaCuO}_4^{+ \pm x}\) has 3D antiferromagnetic ordering of the \(\text{Cu}^{2+}\) spins. Above \(T_N\), the spins are still ordered in two dimensions. As discussed in the last section, two-magnon (spin-pair) scattering occurs as a well-defined peak at about 3000 cm\(^{-1}\) in nonsuperconducting \(\text{La}_2\text{CuO}_4^{0+}\).\(^{20}\) The frequency of one-magnon scattering in \(\text{La}_2\text{CuO}_4^{0+}\) is essentially zero\(^{16}\) near the zone center (i.e., \(k \approx 0\)), as probed by light scattering. Clearly, the 630-cm\(^{-1}\) peak is not associated with two-magnon scattering or one-magnon scattering near the zone center. It is possible that defects such as vacancies could result in observable scattering from a localized, single spin flip. Indeed, a single spin flip next to a vacancy has been suggested as a possible origin of the \(\approx 1430\) cm\(^{-1}\) peak of \(\text{La}_2\text{CuO}_4^{0+}\).\(^{25,28}\) However, a similar explanation is not possible for the 630-cm\(^{-1}\) peak. If one takes the exchange parameter \(J \approx 1000–1300\) cm\(^{-1}\) (Refs. 26 and 29) and \(\text{spin}=\frac{1}{2}\), then 630 cm\(^{-1}\) requires more than two vacancies around a copper atom,\(^{23}\) an unphysical prospect. Since magnetic scattering can be ruled out, it is concluded that the 630-cm\(^{-1}\) peak of SC \(\text{LaCuO}_4^{+ \pm x}\) is phonon scattering.

One interpretation is that the peak at 630 cm\(^{-1}\) is a local phonon mode associated with the excess oxygen in the structure. Following this possibility and literature suggestions that the excess oxygen forms a dioxygen complex in the structure, it is interesting to compare the frequency of \(\text{O}_2\)\(^{-\pm}\) vibration as a function of O—O bond dis-

\begin{table}[h]
\centering
\caption{Vibrational frequency of dioxygen complexes of coordination compounds as a function of valence. Data from Nakamoto (Ref. 31).}
\begin{tabular}{cccccc}
\hline
 & \([\text{O}_2^+]\text{AsF}_6\) & \multicolumn{2}{c}{Molecular} & \multicolumn{2}{c}{Peroxide} & \multicolumn{1}{c}{SC \(\text{LaCuO}_4^{+ \pm x}\)} \\
 & & \begin{tabular}{c}O\(_2\) \end{tabular} & \begin{tabular}{c}K[\text{O}_2^+]\end{tabular} & \begin{tabular}{c}Na\(_2[\text{O}_2\text{O}^\pm]\) \end{tabular} & \begin{tabular}{c}Na\(_2[\text{O}_2\text{O}^\pm]\) \\
\hline
Bond distance (Å) & 1.123 & 1.207 & 1.28 & 1.49 & \(\approx 760\) & 630\(^b\) \\
\(v\) (cm\(^{-1}\)) & 1858 & 1555 & 1108 & & & \\
\end{tabular}
\end{table}

\(^a\)Short O—O bond length as suggested by neutron diffraction (Ref. 5).
\(^b\)Raman peak observed in SC \(\text{LaCuO}_4^{+ \pm x}\) No. 2 but not in NSC \(\text{LaCuO}_4^{+ \pm x}\).
tance. As seen in Table III, the O-O vibration becomes weaker with increasing O—O bond distance. Table III also lists the 630-cm\(^{-1}\) frequency and the short O—O bond length observed by neutron diffraction for La\(_2\)CuO\(_4\), \(D\). While this interpretation cannot be rigorously defended (see below), the relationship would suggest that if the 630-cm\(^{-1}\) peak is indeed from dioxygen vibration, then the dioxygen complex is more like peroxide (O\(_2^{2-}\)) than like superoxide (O\(_2^{-1}\)). The observed polarization dependence of the 630-cm\(^{-1}\) peak is qualitatively consistent with the structural model proposed by Chaloup and co-workers.\(^5,10\) As illustrated in Fig. 1, the short O(3)—O(4) bond of 1.59 Å lies largely within the x-y plane. Therefore, vibration of O(3)—O(4) along this bond should have a large effect on the polarizability measured by (yy) polarization and a smaller influence on the polarizability measured by (zz) polarization. The experimental observations are consistent with this intuition—the 630-cm\(^{-1}\) peak is strong for (yy) polarization and is absent for (zz) polarization. However, arguments based upon strong scattering for (yy) polarization are weakened by the fact that anomalous scattering peaks are observed with (yy) polarization for La\(_2\)CuO\(_4\) (see Sec. V).

Despite the reasoning given above, the 630-cm\(^{-1}\) peak cannot be definitely assigned to vibration of a dioxygen complex. To begin with, 630 cm\(^{-1}\) is within the frequency range expected for vibrational modes of oxygen in the stoichiometric La\(_2\)CuO\(_4\) structure. For example, an X-point mode of tetragonal La\(_2\)CuO\(_4\) has been calculated from first principles to occur at 731 cm\(^{-1}\).\(^30\) This mode is a planar-breathing vibration of O(2) with some contribution from the axial vibration of O(1). In orthorhombic La\(_2\)CuO\(_4\), this mode is folded back into the zone center and is Raman allowed. Indeed, Weber et al.\(^{23}\) have assigned a peak observed at 710 cm\(^{-1}\) in orthorhombic La\(_2\)CuO\(_4\) to the planar-breathing mode of O(2). There exist several mechanisms that might allow phonons that are formally Raman inactive to gain Raman intensity even in tetragonal SC LaCuO\(_{4+x}\). The excess oxygen can be viewed as a defect that destroys the long-range order and local symmetry of the lattice and thus relaxes selection rules associated with momentum conservation and phonon symmetry. This could allow vibrational modes that are not at the zone center or are normally only ir-active to gain Raman intensity. In addition, the SC LaCuO\(_{4+x}\) structure may have a larger unit cell than even the underlying orthorhombic structure of La\(_2\)CuO\(_4\). This large unit cell could have additional Raman-allowed phonons not directly related to the speculated dioxygen species and the 630-cm\(^{-1}\) peak could be such a phonon. Finally, the 630-cm\(^{-1}\) peak could be a second-order phonon.

Despite the complications discussed above, several conclusions can be drawn. The 630-cm\(^{-1}\) peak of SC LaCuO\(_{4+x}\) No. 2 is of phonon origin. Since it is not present in NSC LaCuO\(_{4.0}\) or La\(_{2-x}\)Sr\(_x\)CuO\(_4\),\(^6\) the 630-cm\(^{-1}\) peak results from the introduction of excess oxygen into the structure. Further, the 630-cm\(^{-1}\) peak is not related to either the orthorhombic distortion or phase separation since it occurs in tetragonal SC LaCuO\(_{4+x}\).

VII. SUMMARY

In certain aspects, the excess oxygen of SC LaCuO\(_{4+x}\) has little influence on the Raman-allowed phonons. For example, at low temperature, five \(A_g\) phonons are found at essentially the same frequency in both SC LaCuO\(_{4+x}\) and NSC LaCuO\(_{4.0}\). Further, like NSC LaCuO\(_{4.0}\), SC LaCuO\(_{4+x}\) has a soft-mode phonon associated with the orthorhombic-to-tetragonal transition and the temperature dependence of this mode below about 250 K is nearly identical in the two materials. However, key differences do exist between SC LaCuO\(_{4+x}\) and NSC LaCuO\(_{4.0}\). The soft \(A_g\) phonon of SC LaCuO\(_{4+x}\) is only observed below about 250 K, approximately the temperature of phase transition and phase separation.\(^5,12,14\) In contrast, the T-O transition in NSC LaCuO\(_{4.0}\) occurs at about 530 K. At room temperature, SC LaCuO\(_{4+x}\) does not have a well-defined peak from two-magnon scattering, in contrast to NSC LaCuO\(_{4.0}\). At 77 K (i.e., in its phase-separated form), SC LaCuO\(_{4+x}\) exhibits a well-defined peak from two-magnon scattering. This establishes that the La\(_2\)CuO\(_4\) present in phase separated SC LaCuO\(_{4+x}\) is antiferromagnetic and not superconducting. In both the tetragonal and orthorhombic forms, certain regimes of SC LaCuO\(_{4+x}\) have a phonon peak at 630 cm\(^{-1}\), unlike NSC LaCuO\(_{4.0}\). This peak is believed to occur in samples containing the most excess oxygen. If this peak results from a local vibrational mode of a dioxygen species in the structure, then the observed frequency is consistent with peroxide but not superoxide. However, a definitive assignment of the 630-cm\(^{-1}\) phonon is not possible at this time.

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