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Authors
Atlas, C
Aad, G
Abbott, B
et al.

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Measurement of the charged-particle multiplicity inside jets from √s = 8 TeV pp collisions with the ATLAS detector

ATLAS Collaboration

CERN, 1211 Geneva 23, Switzerland

Abstract The number of charged particles inside jets is a widely used discriminant for identifying the quark or gluon nature of the initiating parton and is sensitive to both the perturbative and non-perturbative components of fragmentation. This paper presents a measurement of the average number of charged particles with p_T > 500 MeV inside high-momentum jets in dijet events using 20.3 fb⁻¹ of data recorded with the ATLAS detector in pp collisions at √s = 8 TeV collisions at the LHC. The jets considered have transverse momenta from 50 GeV up to and beyond 1.5 TeV. The reconstructed charged-particle track multiplicity distribution is unfolded to remove distortions from detector effects and the resulting charged-particle multiplicity is compared to several models. Furthermore, quark and gluon jet fractions are used to extract the average charged-particle multiplicity for quark and gluon jets separately.

1 Introduction

Quarks and gluons produced in high-energy particle collisions hadronize before they can be observed directly. However, the properties of the resulting collimated sprays of hadrons, known as jets, depend on the type of parton which initiated them. One jet observable sensitive to the quark or gluon nature is the number of charged particles inside the jet. Due to their larger colour-charge under the strong force, gluon-initiated jets contain on average more particles than quark-initiated jets. The average (charged) particle multiplicity inside jets increases with jet energy, but increases faster for gluon-initiated jets than for quark-initiated jets [1].

These properties were used recently at the Large Hadron Collider (LHC) to differentiate between jets originating from a quark or a gluon [2–6]. These studies have found significant differences in the charged-particle multiplicity between the available simulations and data. Improved modelling based on measurements of the number of charged particles inside jets is thus crucial for future studies.

This paper presents a measurement of the average charged-particle multiplicity inside jets as a function of the jet transverse momentum in dijet events in pp collisions at √s = 8 TeV with the ATLAS detector. The measurement of the charged-particle multiplicity inside jets has a long history from the SPS [7–9], PETRA [10,11], PEP [12–15], TRISTAN [16], CESR [17], LEP [18–29], and the Tevatron [30]. At the LHC, both ATLAS [31,32] and CMS [33] have measured the charged-particle multiplicity inside jets at √s = 7 TeV. One ATLAS result [31] used jets that are reconstructed using tracks and have transverse momentum less than 40 GeV. A second ATLAS analysis [32] has measured charged particles inside jets with transverse momenta spanning the range from 50 to 500 GeV with approximately constant 3–4 % uncertainties. The CMS measurement [33] spans jet transverse momenta between 50 and 800 GeV with 5–10 % uncertainties in the bins of highest transverse momentum. The analysis presented here uses the full √s = 8 TeV ATLAS dataset, which allows for a significant improvement in the precision at high transverse momentum up to and beyond 1.5 TeV.

This paper is organized as follows. After a description of the ATLAS detector and object and event selection in Sect. 2, simulated samples are described in Sect. 3. In order for the measured charged-particle multiplicity to be compared with particle-level models, the data are unfolded to remove distortions from detector effects, as described in Sect. 4. Systematic uncertainties in the measured charged-particle multiplicity are discussed in Sect. 5 and the results are presented in Sect. 6.

2 Object and event selection

ATLAS is a general-purpose detector designed to measure the properties of particles produced in high-energy pp collisions with nearly a full 4π coverage in solid angle. Charged-
particle momenta are measured by a series of tracking detectors covering a range of $|\eta| < 2.5$ and immersed in a 2 T axial magnetic field, providing measurements of the transverse momentum, $p_T$, with a resolution $\sigma_{p_T}/p_T \sim 0.05 \% \times p_T/\text{GeV} \pm 1 \%$. Electromagnetic and hadronic calorimeters surround the tracking detector, with forward calorimeters allowing electromagnetic and hadronic energy measurements up to $|\eta| = 4.5$. A detailed description of the ATLAS detector can be found in Ref. [34].

This measurement uses the dataset of $pp$ collisions recorded by the ATLAS detector in 2012, corresponding to an integrated luminosity of 20.3 fb$^{-1}$ at a center-of-mass energy of $\sqrt{s} = 8 \text{ TeV}$. The data acquisition and object/event selection are described in detail in Ref. [35] and highlighted here for completeness. Jets are clustered using the anti-$k_T$ jet algorithm [36] with radius parameter $R = 0.4$ implemented in FastJet [37] using as inputs topological calorimeter-cell clusters [38], calibrated using the local cluster weighting (LCW) algorithm [39,40]. An overall jet energy calibration accounts for residual detector effects as well as contributions from multiple proton–proton collisions in the same bunch crossing (pileup) [41] in order to make the reconstructed jet energy correspond to an unbiased measurement of the particle-level jet energy. Jets are required to be central ($|\eta| < 2.1$) so that their charged particles are within the $|\eta| < 1.5$ coverage of the tracking detector. Events are further required to have at least two jets with $p_T > 50 \text{ GeV}$ and only the leading two jets are considered for the charged-particle multiplicity measurement. To select dijet topologies where the jets are balanced in $p_T$, the two leading jets must have $p_T^\text{lead}/p_T^\text{sublead} < 1.5$, where $p_T^\text{lead}$ and $p_T^\text{sublead}$ are the transverse momenta of the jets with the highest and second-highest $p_T$, respectively. The jet with the smaller (larger) absolute pseudorapidity $|\eta|$ is classified as the more central (more forward) jet. A measurement of the more forward and more central average charged-particle multiplicities can exploit the rapidity dependence of the jet type to extract information about the multiplicity for quark- and gluon-initiated jets as is described in Sect. 6. The more forward jet tends to be correlated with the parton with higher longitudinal momentum fraction $x$, and is less likely to be a gluon-initiated jet.

Tracks are required to have $p_T \geq 500 \text{ MeV}$, $|\eta| < 2.5$, and a $\chi^2$ per degree of freedom (resulting from the track fit) less than 3.0. Additional quality criteria are applied to select tracks originating from the collision vertex and reject fake tracks reconstructed from random hits in the detector. In particular, tracks are matched to the hard-scatter vertex by requiring $|z_0 \sin(\theta)| < 1.5 \text{ mm}$ and $|d_0| < 1 \text{ mm}$, where $z_0$ and $d_0$ are the track longitudinal and transverse impact parameters, respectively, calculated with respect to the primary vertex. Tracks must furthermore have at least one hit in the silicon pixel detector and at least six hits in the semiconductor microstrip detector.

The matching of tracks with the calorimeter-based jets is performed via the ghost-association technique [42]: the jet clustering process is repeated with the parton with higher longitudinal momentum fraction $x$, and is less likely to be a gluon-initiated jet.

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3 Event simulation

Monte Carlo (MC) samples are used in order to determine how the detector response affects the charged-particle multiplicity and to make comparisons with the corrected data. The details of the sample used are shown in Table 1. The sample is generated with Pythia 8.1.75 [43] using the AU2 [44] set of tuned parameters (tune) and the HERWIG++ 2.6.3 [45]
sample with the UE-EE [46] tune are further processed with the ATLAS detector simulation [47] based on GEANT4 [48]. The effects of pileup are modelled by adding to the generated hard-scatter events (before the detector simulation) multiple minimum-bias events generated with Pythia 8.160, the A2 tune [44], and the MSTW2008LO [49] Parton distribution function (PDF) set. The distribution of the number of interactions is then weighted to reflect the pileup distribution in the data.

4 Unfolding

The measurement is carried out within a fiducial volume matching the experimental selection to avoid extrapolation into unmeasured kinematic regions that have additional model dependence and related uncertainties. Particle-level definitions of the reconstructed objects are chosen to be as close as possible to those described in Sect. 2. Particle-level jets are clustered from generated stable particles with a mean lifetime $\tau > 30$ ps, excluding muons and neutrinos. As with the detector-level jets, particle-level jets are clustered with the anti-$k_t$, $R = 0.4$ algorithm. Any charged particle clustered in a particle-level jet is considered for the charged-particle multiplicity calculation if it has $p_T > 500$ MeV. Events are required to have at least two jets with $|\eta| < 2.1$ and $p_T > 50$ GeV and the two highest-$p_T$ jets must satisfy the same $p_T$-balance requirement between the leading and sub-leading jet as at detector level ($p_{T,\text{lead}}^1/p_{T,\text{sublead}}^1 < 1.5$). The $p_T$ symmetry requirement enriches the sample in a back-to-back topology and suppresses non-isolated jets. In more than 70% of events, the nearest jet in $\Delta R$ with $p_T > 25$ GeV is the other selected jet and in less than 7% of events, there is a jet with $p_T > 25$ GeV within $\Delta R = 0.8$ from one of the two selected jets. Due to the high-energy and well-separated nature of the selected jets, the hard-scatter quarks and gluons can be cleanly matched to the outgoing jets. In this analysis, the type of a jet is defined as that of the highest-energy parton in simulation within a $\Delta R = 0.4$ cone around the particle-jet’s axis. Figure 2 shows the fraction of gluon-initiated jets as a function of jet $p_T$ for the more forward and more central jet within the event. The fraction of gluon-initiated jets decreases with $p_T$, but the difference between the more forward and more central jets peaks around $p_T \sim 350$ GeV. This difference is exploited in Sect. 6 to extract separately the average quark- and gluon-initiated jet charged-particle multiplicity.

The average charged-particle multiplicity in particle-level jets is determined as a function of jet $p_T$. An iterative Bayesian (IB) technique [61] as implemented in the RooUnfold framework [62] is used to unfold the two-dimensional charged-particle multiplicity and jet $p_T$ distribution. In the IB unfolding technique, the number of iterations and the prior distribution are the input parameters. The raw data are corrected using the simulation to account for events that pass the fiducial selection at detector level, but not the corresponding selection at particle level; this correction is the fake factor. Then, the IB method iteratively applies Bayes’ theorem using the response matrix to con-

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**Table 1** Monte Carlo samples used in this analysis. The abbreviations ME, PDF, and UE respectively stand for matrix element, parton distribution function, and underlying event. “Tune” refers to the set of tunable MC parameters used.

<table>
<thead>
<tr>
<th>ME generator</th>
<th>PDF</th>
<th>Tune</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pythia 8.175 [43]</td>
<td>CT10 [50]</td>
<td>AU2 [44]</td>
</tr>
<tr>
<td>Pythia 8.186</td>
<td>NNPDF2.3 [51]</td>
<td>Monash [52]</td>
</tr>
<tr>
<td>Pythia 8.186</td>
<td>NNPDF2.3</td>
<td>A14 [53]</td>
</tr>
<tr>
<td>Herwig++ 2.6.3</td>
<td>CTEQ6L1</td>
<td>UE-EE3 [46]</td>
</tr>
<tr>
<td>Herwig++ 2.7.1</td>
<td>CTEQ6L1</td>
<td>UE-EE5 [57]</td>
</tr>
<tr>
<td>Pythia 6.428 [58]</td>
<td>CTEQ6L1</td>
<td>P2012 [59]</td>
</tr>
<tr>
<td>Pythia 6.428</td>
<td>CTEQ6L1</td>
<td>P2012RadLo [59]</td>
</tr>
<tr>
<td>Pythia 6.428</td>
<td>CTEQ6L1</td>
<td>P2012RadHi [59]</td>
</tr>
</tbody>
</table>

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**Fig. 2** The simulated fraction of jets originating from gluons as a function of jet $p_T$ for the more forward jet (down triangle), the more central jet (up triangle), and the difference between these two fractions (circle). The fractions are derived from Pythia 8 with the CT10 PDF set and the error bars represent the PDF and matrix element uncertainties, further discussed in Sect. 6. The uncertainties on the fraction difference are computed from propagating the uncertainties on the more forward and more central fractions, treating as fully correlated.

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While it is possible to classify jets as quark- or gluon-initiated beyond leading order in $m_{jet}/E_{jet}$ [60], the classification is jet algorithm-dependent and unnecessary for the present considerations. For the results presented in Sect. 6 that rely on jet-type labelling, alternative definitions were considered and found to have a negligible impact compared to other sources of theoretical and experimental uncertainty.
Fig. 3 The jet $p_T$ dependence of a the average reconstructed track multiplicity for uncorrected data and detector-level simulation, b the average reconstructed track multiplicity for the detector-level simulation and the average charged-particle multiplicity for the particle-level simulation, c the average charged-particle multiplicity for the unfolded data and the particle-level simulation, and d the average charged-particle multiplicity divided by the average reconstructed track multiplicity in simulation. Three charged-particle and track $p_T$ thresholds are used in each case: 0.5, 2, and 5 GeV. 

Pythia 8 with the CT10 PDF and the AU2 tune are used for the simulation. For the data, only statistical uncertainties are included in the error bars (which are smaller than the markers for most bins).

The number of iterations in the IB method trades off unfolding bias against statistical fluctuations. An optimal value of four iterations is obtained by minimizing the bias when unfolding pseudo-data derived from Herwig++ using a prior distribution and a response matrix derived from Pythia as a test of the methodology. Lastly, unfolding applies another correction from simulation to the unfolded data to account for events passing the particle-level selection but not the detector-level selection; this correction is the inefficiency factor.

Figure 3 displays the $p_T$ dependence of the average charged-particle multiplicity for uncorrected data and detector-level simulation and for particle-level simulation as well as the unfolded data. The prediction from Pythia 8 with the AU2 tune has too many tracks compared with the uncorrected data, and the size of the data/MC difference increases with decreasing track $p_T$ threshold (Fig. 3a). The difference between the detector-level and particle-level simulation in...
Fig. 3b (for which the ratio is given in Fig. 3d) gives an indication of the corrections required to account for detector acceptance and resolution effects in the unfolding procedure. Particle-level distributions in Fig. 3c show similar trends to the detector-level ones in Fig. 3a.

5 Systematic uncertainties

All stages of the charged-particle multiplicity measurement are sensitive to sources of potential bias. The three stages of the measurement are listed below, with an overview of the systematic uncertainties that impact the results at each stage:

Response matrix For events that pass both the detector-level and particle-level fiducial selections, the response matrix describes migrations between bins when moving between the detector level and the particle level. The response matrix is taken from simulation and various experimental uncertainties in the charged-particle multiplicity and jet $p_T$ spectra result in uncertainties in the matrix. These uncertainties can be divided into two classes: those impacting the calorimeter-based jet $p_T$ and those impacting track reconstruction inside jets. The dominant uncertainty at high jet $p_T$ is due to the loss of charged-particle tracks in the jet core due to track merging. This charged energy loss uncertainty is estimated using the data/MC differences in the ratio of the track-based jet $p_T$ to the calorimeter-based jet $p_T$ [35]. More charged energy is lost in the data than in the MC and thus this uncertainty is one-sided. There are other tracking uncertainties in the track momentum scale and resolution, the track reconstruction efficiency, and the rate of tracks formed from random combinations of hits (fake tracks). The prescription for these sub-dominant tracking uncertainties is identical to Ref. [35]. The uncertainties related to the calorimeter-based jet are sub-dominant (except in the lowest $p_T$ bins) and are due to the uncertainty in the jet energy scale and the jet energy resolution.

Correction factors Fake and inefficiency factors are derived from simulation to account for the fraction of events that pass either the detector-level or particle-level fiducial selection ($p_T > 50$ GeV $|\eta| < 2.1$, and $P_{lead}^T/P_{sublead}^T < 1.5$), but not both. These factors are derived in bins of jet $p_T$ and charged particle multiplicity, separately for the more forward and more central jets. They are generally higher than the data and this is more pronounced at higher jet $p_T$. The default ATLAS tune in Run 1 (AU2) performs similarly to the Monash tune, but the prediction with A14 (the ATLAS default for the analysis of Run 2 data) is significantly closer to the data. A previous ATLAS measurement [31] of charged-particle multiplicity inside jets was included in the tuning of A14, but the jets in that measurement have $p_T \lesssim 50$ GeV. One important difference between A14 and Monash is that the value of $\alpha_s$ governing the amount of final-state radiation is about 10% lower in A14 than in Monash. This parameter has a large impact on the average charged-particle multiplicity, which is shown by the PYTHIA 6 lines in Fig. 4 where the Perugia radHi and radLo tunes are significantly separated from the central P2012 tune. The $\alpha_s$ value that regulates final-state radiation is changed by factors of one half and two for these tunes with respect to the nominal Perugia 2012 tune. The recent (and Run 2 default) EE5 underlying-event tune for HERWIG++ improves the modelling of the average charged-particle multiplicity with respect to the EE3 tune (Run 1 default).

6 Results

The unfolded average charged-particle multiplicity combining both the more forward and the more central jets is shown in Fig. 4, compared with various model predictions. As was already observed for the reconstructed data in Fig. 1, the average charged-particle multiplicity in data falls between the predictions of PYTHIA 8 and HERWIG++, independently of the underlying-event tunes. The PYTHIA 8 predictions are generally higher than the data and this is more pronounced at higher jet $p_T$. The default ATLAS tune in Run 1 (AU2) performs similarly to the Monash tune, but the prediction with A14 (the ATLAS default for the analysis of Run 2 data) is significantly closer to the data. A previous ATLAS measurement [31] of charged-particle multiplicity inside jets was included in the tuning of A14, but the jets in that measurement have $p_T \lesssim 50$ GeV. One important difference between A14 and Monash is that the value of $\alpha_s$ governing the amount of final-state radiation is about 10% lower in A14 than in Monash. This parameter has a large impact on the average charged-particle multiplicity, which is shown by the PYTHIA 6 lines in Fig. 4 where the Perugia radHi and radLo tunes are significantly separated from the central P2012 tune. The $\alpha_s$ value that regulates final-state radiation is changed by factors of one half and two for these tunes with respect to the nominal Perugia 2012 tune. The recent (and Run 2 default) EE5 underlying-event tune for HERWIG++ improves the modelling of the average charged-particle multiplicity with respect to the EE3 tune (Run 1 default).
A summary of all the systematic uncertainties and their impact on the charged-particle multiplicity $n_{\text{charged}}$ in the $p_T$ range $0.5 < p_T < 5$ GeV, and the more central jet. Uncertainties are given in percent. A value of 0.0 is quoted if the uncertainty is below 0.5%.

The difference in the average charged-particle multiplicity between the more forward and the more central jet is sensitive to the difference between quark and gluon constituent multiplicities. Figure 5a shows that the difference is significant for $p_T \lesssim 1.1$ TeV. The shape is governed by the difference in the gluon fraction between the more forward and the more central jet, which was shown in Fig. 2 to peak around $p_T \sim 350$ GeV. The average difference, combined with the gluon fraction, can be used to extract the average charged-particle multiplicity for quark- and gluon-initiated jets separately. Given the quark and gluon fractions $f_q(x)$ and $f_g(x)$, the average charged-particle multiplicity for quark- and gluon-initiated jets is extracted by solving the system of equations in Eq. (1):

$$
\langle n_{\text{charged}}^f \rangle = f_q^f \langle n_{\text{charged}}^q \rangle + f_g^f \langle n_{\text{charged}}^g \rangle
$$

Given the jet $p_T$, the charged particle multiplicity inside jets does not vary significantly with $\eta$. This is confirmed by checking that the solution to Eq. 1 reproduces the quark and gluon jet charged particle multiplicities for both PYTHIA 8 and HERWIG++ to better than 1% across most of the $p_T$ range. The extracted $p_T$ dependence of the average charged-particle multiplicities for quark- and gluon-initiated jets is shown in Fig. 5b. PYTHIA 8 with the CT10 PDF set is used to determine the gluon fractions. The experimental uncertainties are propagated through Eq. (1) by recomputing the quark and gluon average charged-particle multiplicities for each variation accounting for a systematic uncertainty; the more forward and more central jet uncertainties are treated as being fully correlated. In addition to the experimental uncertainties, the error bands in Fig. 5b include uncertainties in the gluon fractions from both the PDF and matrix element (ME) uncertainties. The PDF uncertainty is determined using the CT10 eigenvector PDF sets and validated by comparing CT10 and NNPDF. The ME uncertainty is estimated by comparing the fractions $f_q^{f/c}$ from PYTHIA 8 and HERWIG++ after reweighting the PYTHIA 8 sample with CT10 to CTEQ6L1 to match the PDF used for HERWIG++. All PDF re-weighting is performed using LHAPDF6 [64]. The PDF and ME uncertainties are comparable in size to the total experimental uncertainty. As expected, the average multiplicity increases with jet $p_T$ for both the quark-initiated jets and gluon-initiated jets, but increases faster for gluon-initiated jets. Furthermore, the multiplicity is significantly higher for gluon-initiated jets than for quark-initiated jets. The average charged-particle multiplicity in PYTHIA 8 with the AU2 tune is higher than in the data for both the quark- and gluon-initiated jets. In addition to predictions from leading-logarithm parton shower simulations, calcu-
Fig. 4 The measured average charged-particle multiplicity as a function of the jet $p_T$, combining the more forward and the more central jets for (a) $p_T^{\text{track}} > 0.5$ GeV, (b) $p_T^{\text{track}} > 2$ GeV, and (c) $p_T^{\text{track}} > 5$ GeV. The band around the data is the sum in quadrature of the statistical and systematic uncertainties. Error bars on the data points represent the statistical uncertainty (which are smaller than the markers for most bins).

Simulations of the scale dependence for the parton multiplicity inside jets have been performed in perturbative quantum chromodynamics (pQCD). Up to a non-perturbative factor that is constant for the jet $p_T$ range considered in this analysis, these calculations can be interpreted as a prediction for the scale dependence of $\langle n_{\text{charged}} \rangle$ for quark- and gluon-initiated jets. There are further caveats to the predictability of such a calculation since $n_{\text{charged}}$ is not infrared safe or even Sudakov safe [65]. Therefore, the formal accuracy of the series expansion in $\sqrt{\alpha_s}$ is unknown. Given these caveats, the next-to-next-to-next-to-leading-order (N$^3$LO) pQCD calculation [66,67] is overlaid in Fig. 5 with renormalization scale $\mu = R p_T$ in the five-flavour scheme and $R = 0.4$. The theoretical error band is calculated by varying $\mu$ by a factor of two. The prediction cannot give the absolute scale,

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This factor is found to be about 0.19 for gluon jets and 0.25 for quark-initiated jets.
and therefore the curve is normalized in the second $p_T$ bin (100 GeV < $p_T$ < 200 GeV) where the statistical uncertainty is small. The predicted scale dependence for gluon-initiated jets is consistent with the data within the uncertainty bands while the curve for quark-initiated jets is higher than the data by about one standard deviation.

7 Summary

This paper presents a measurement of the $p_T$ dependence of the average jet charged-particle multiplicity in dijet events from 20.3 fb$^{-1}$ of $\sqrt{s}$ = 8 TeV $pp$ collision data recorded by the ATLAS detector at the LHC. The measured charged-particle multiplicity distribution is unfolded to correct for the detector acceptance and resolution to facilitate direct comparison to particle-level models. Comparisons are made at particle level between the measured average charged-particle multiplicity and various models of jet formation. Significant differences are observed between the simulations using Run 1 tunes and the data, but the Run 2 tunes for both PYTHIA 8 and HERWIG++ significantly improve the modelling of the average $n_{\text{charge}}$. Furthermore, quark- and gluon-initiated jet constituent charged-particle multiplicities are extracted and compared with simulations and calculations. As expected, the extracted gluon-initiated jet constituent charged-particle multiplicity is higher than the corresponding quantity for quark-initiated jets and a calculation of the $p_T$-dependence accurately models the trend observed in the data.

The particle-level spectra are available [68] for further interpretation and can serve as a benchmark for future measurements of the evolution of non-perturbative jet observables to validate MC predictions and tune their model parameters.

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W. K. Balunas123, E. Banas40, Sw. Banerjee172,e, A. A. E. Bannoura174, L. Barak31, E. L. Barberio90, D. Barberis31a,51b,

1 Department of Physics, University of Adelaide, Adelaide, Australia
2 Physics Department, SUNY Albany, Albany, NY, USA
3 Department of Physics, University of Alberta, Edmonton, AB, Canada
4 (a)Department of Physics, Ankara University, Ankara, Turkey; (b)Istanbul Aydin University, Istanbul, Turkey; (c)Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey
5 LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France
6 High Energy Physics Division, Argonne National Laboratory, Argonne, IL, USA
7 Department of Physics, University of Arizona, Tucson, AZ, USA
8 Department of Physics, The University of Texas at Arlington, Arlington, TX, USA
9 Physics Department, University of Athens, Athens, Greece
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
12 Institut de Física d’Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain
13 Institute of Physics, University of Belgrade, Belgrade, Serbia
14 Department for Physics and Technology, University of Bergen, Bergen, Norway
15 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, USA
16 Department of Physics, Humboldt University, Berlin, Germany
17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, UK
19 (a)Department of Physics, Bogazici University, Istanbul, Turkey; (b)Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey; (c)Faculty of Engineering and Natural Sciences, Istanbul Bilgi University, Istanbul, Turkey; (d)Faculty of Engineering and Natural Sciences, Bahcesehir University, Istanbul, Turkey
20 Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
21 (a)INFN Sezione di Bologna, Bologna, Italy; (b)Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
22 Physikalisches Institut, University of Bonn, Bonn, Germany
23 Department of Physics, Boston University, Boston, MA, USA
24 Department of Physics, Brandeis University, Waltham, MA, USA
25 (a)Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil; (b)Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil; (c)Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil; (d)Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
26 Physics Department, Brookhaven National Laboratory, Upton, NY, USA
27 (a)Transilvania University of Brasov, Brasov, Romania; (b)National Institute of Physics and Nuclear Engineering, Bucharest, Romania; (c)Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj Napoca, Romania; (d)University Politehnica Bucharest, Bucharest, Romania; (e)West University in Timisoara, Timisoara, Romania
28 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
29 Cavendish Laboratory, University of Cambridge, Cambridge, UK
30 Department of Physics, Carleton University, Ottawa, ON, Canada
31 CERN, Geneva, Switzerland
32 Enrico Fermi Institute, University of Chicago, Chicago, IL, USA
33 (a)Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile; (b)Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
34 (a)Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China; (b)Department of Modern Physics, University of Science and Technology of China, Hefei, Anhui, China; (c)Department of Physics, Nanjing University, Nanjing, Jiangsu, China; (d)School of Physics, Shandong University, Jinan, Shandong, China; (e)Shanghai Key Laboratory for Particle Physics and Cosmology, Department of Physics and Astronomy, Shanghai Jiao Tong University, also affiliated with PKU-CHEP, Shanghai, China; (f)Physics Department, Tsinghua University, Beijing 100084, China

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35 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
36 Nevis Laboratory, Columbia University, Irvington, NY, USA
37 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
38 (a) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Frascati, Italy; (b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
39 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Kraków, Poland; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Kraków, Poland
40 Institute of Nuclear Physics, Polish Academy of Sciences, Kraków, Poland
41 Physics Department, Southern Methodist University, Dallas, TX, USA
42 Physics Department, University of Texas at Dallas, Richardson, TX, USA
43 DESY, Hamburg and Zeuthen, Germany
44 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
45 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
46 Department of Physics, Duke University, Durham, NC, USA
47 SUPA-School of Physics and Astronomy, University of Edinburgh, Edinburgh, UK
48 INFN Laboratori Nazionali di Frascati, Frascati, Italy
49 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
50 Section de Physique, Université de Genève, Geneva, Switzerland
51 (a) INFN Sezione di Genova, Genoa, Italy; (b) Dipartimento di Fisica, Università di Genova, Genoa, Italy
52 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
53 II Physikalisches Institut, Justus-Liebig-Universität Gießen, Gießen, Germany
54 SUPA-School of Physics and Astronomy, University of Glasgow, Glasgow, UK
55 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
56 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
57 Department of Physics, Hampton University, Hampton, VA, USA
58 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, USA
59 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany; (c) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
60 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
61 (a) Department of Physics, The Chinese University of Hong Kong, Shatin, NT, Hong Kong; (b) Department of Physics, The University of Hong Kong, Hong Kong, China; (c) Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
62 Department of Physics, Indiana University, Bloomington, IN, USA
63 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
64 University of Iowa, Iowa City, IA, USA
65 Department of Physics and Astronomy, Iowa State University, Ames, IA, USA
66 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
67 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
68 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
69 Graduate School of Science, Kobe University, Kobe, Japan
70 Faculty of Science, Kyoto University, Kyoto, Japan
71 Kyoto University of Education, Kyoto, Japan
72 Department of Physics, Kyushu University, Fukuoka, Japan
73 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
74 Physics Department, Lancaster University, Lancaster, UK
75 (a) INFN Sezione di Lecce, Lecce, Italy; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
76 Oliver Lodge Laboratory, University of Liverpool, Liverpool, UK
77 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
78 School of Physics and Astronomy, Queen Mary University of London, London, UK
79 Department of Physics, Royal Holloway University of London, Surrey, UK
Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver, BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada

Department of Physics, University of Warwick, Coventry, UK

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, WI, USA

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, CT, USA

Yerevan Physics Institute, Yerevan, Armenia

Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

d Also at Department of Physics, King’s College London, London, UK
b Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
c Also at Novosibirsk State University, Novosibirsk, Russia
d Also at TRIUMF, Vancouver, BC, Canada
e Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY, USA
f Also at Department of Physics, California State University, Fresno, CA, USA
g Also at Department of Physics, University of Fribourg, Fribourg, Switzerland

h Also at Departamento de Física de la Universitat Autonoma de Barcelona, Barcelona, Spain
i Also at Departamento de Física e Astronomia, Faculdade de Ciencias, Universidade do Porto, Porto, Portugal
j Also at Tomsk State University, Tomsk, Russia
k Also at Universita di Napoli Parthenope, Naples, Italy
l Also at Institute of Particle Physics (IPP), Canada

m Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia
n Also at Department of Physics, The University of Michigan, Ann Arbor, MI, USA
o Also at Louisiana Tech University, Ruston, LA, USA
p Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain
q Also at Graduate School of Science, Osaka University, Osaka, Japan
r Also at Department of Physics, National Tsing Hua University, Hsinchu City, Taiwan
s Also at Department of Physics, The University of Texas at Austin, Austin, TX, USA
t Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia
u Also at CERN, Geneva, Switzerland
v Also at Georgian Technical University (GTU), Tbilisi, Georgia
w Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan
x Also at Manhattan College, New York, NY, USA
y Also at Hellenic Open University, Patras, Greece
z Also at Institute of Physics, Academia Sinica, Taipei, Taiwan

aa Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
ab Also at School of Physics, Shandong University, Shandong, China
ac Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
ad Also at Section de Physique, Université de Genève, Geneva, Switzerland
ae Also at International School for Advanced Studies (SISSA), Trieste, Italy
af Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, USA
ag Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China
ah Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria
ai Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
aj Also at National Research Nuclear University MEPhI, Moscow, Russia
ak Also at Department of Physics, Stanford University, Stanford, CA, USA
al Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
am Also at Flensburg University of Applied Sciences, Flensburg, Germany
an Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia
ao Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
∗ Deceased