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Aperiodic-Tiling-Based Mushroom-Type High-Impedance Surfaces

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Abstract—This letter deals with a study of aperiodically ordered textured (mushroom-type) high-impedance surfaces (HISs), and their possible application as artificial-magnetic-conductor ground-planes for low-profile directive antennas. In this framework, results from full-wave simulations are presented in order to characterize the electromagnetic response (in terms of impedance matching, directivity, and surface-wave suppression) of selected configurations based on representative categories of aperiodic tilings, and compare them with those of standard periodic HISs.

Index Terms—Antenna ground-planes, artificial impedance surfaces, artificial magnetic conductors (AMCs), quasi-crystals.

I. INTRODUCTION AND BACKGROUND

During the last decade, artificial high-impedance surfaces (HISs) have been extensively investigated and proposed as effective ground planes for improving the electrical and radiative responses of low-profile planar antennas (see [1] for a collection of recent results and applications). The main interest originates from their potential capabilities, within certain frequency ranges, of reflecting electromagnetic (EM) waves with no phase reversal, i.e., behaving as an artificial magnetic conductor (AMC), and suppressing the surface-wave (SW) propagation [2].

Nearly all available examples of HIS configurations are based on particular textures, such as corrugations, mushroom-type protrusions, etc., arranged periodically, for which the EM response is well understood (see, e.g., [3] and [4]). However, inspired by the discovery in solid-state physics of “quasi-crystals” [5], [6] and by the theory of “aperiodic tilings” [7], it has been shown in other artificial-material contexts (e.g., photonic crystals, arrays of subwavelength holes, etc.) that spatial periodicity is not essential to achieve anomalous EM effects like bandgaps, negative refraction and superlensing, directive emission, enhanced transmission, etc. (see [8]–[17] and the references therein for a sparse sampling). In fact, judicious exploitation of the additional degrees of freedom and higher (local, statistical) rotational symmetries potentially available in aperiodically ordered geometries can be exploited to achieve larger bandgaps, lower and/or multiple frequencies of operation, higher isotropy, richer and more frequency-selective defect states, etc. In this framework, a novel mushroom-type HIS configuration was recently proposed in [18], based on a quasi-periodic octagonal (Ammann-Beenker) tiling geometry [7]. Experimental and numerical studies of this structure revealed a very interesting EM response, in terms of multiple frequencies of operation, highly directive radiation from a small electric dipole laid on it, and broadband suppression of the transverse-electric SW.

In this letter, in order to gain further insight into the underlying phenomenology and potentials, we address a comparative study of various HIS configurations with aperiodically ordered mushroom-type textures, based on representative categories of aperiodic tilings. Our study is principally aimed at exploring to what extent the intriguing results observed in [18] are restricted to that particular (octagonal) geometry, or instead they are representative of intrinsic properties of rather generic aperiodically ordered configurations. In this framework, we study three aperiodic representative finite-size HIS ground-planes, as well as the well-known periodic (square) counterpart. Via full-wave simulations, we compare their responses in terms of the radiation (return loss and directivity) from a small electric dipole laid on them, and SW attenuation.

II. GEOMETRY AND PARAMETERS

The three representative aperiodic tilings considered are as follows.

- **Octagonal** (Ammann-Beenker), as in [18], with square- and rhombus-shaped tiles, and eightfold symmetry [7].
- **Dodecagonal**, with square- and equilater-triangle-shaped tiles, and 12-fold symmetry [19].
- **Penrose**, with thick/thin-rhombus-shaped tiles, and fivefold symmetry [7].

The octagonal and Penrose tilings are generated via “cut-and-projection” algorithms [7], whereas the dodecagonal tiling is generated using the Stampfli iterative construction [19]. As a reference configuration, we also consider a standard periodic (square) geometry, whose EM response is well known [2]–[4]. Note that standard concepts and tools typically utilized for the study of periodic HIS configurations (“unit cell,” band-structure, Brillouin zone, etc.) cannot be applied to the aperiodic scenarios of interest here. Therefore, our study involves finite-size HIS structures, obtained by cutting suitably sized regions...
of the above tilings (preserving the local symmetry centers), and placing metallic (copper) patches shaped according to the tiles on top of a 1.6-mm-thick dielectric substrate with relative permittivity \( \varepsilon_r = 2.2 \) (RT/duriod 5880) backed by a metallic ground plane. The patches are suitably scaled, so as to guarantee a constant 0.7-mm gap between them, and are connected to the ground plane by metallic vias of diameter 0.7 mm. The patch characteristic sizes are chosen as in [18], so as to facilitate direct comparison with our results (at least for the octagonal geometry), whereas the total structure sizes are set by our current computational resource limitations. For the various geometries, these values were tuned so as to maintain the same spacing between the metal patches and a comparable number of patches.

Fig. 1(a)–(d) shows the top view of the four HIS configurations under analysis. For the octagonal case, a three-dimensional view is also displayed in Fig. 1(e).

Fig. 1. HIS geometry and parameters. (a)–(d) Top view. (e) Three-dimensional view (octagonal). Metallic (copper) patches of sidelength \( a_{\text{patch}} \) are laid with a 0.7-mm gap on a metal-backed dielectric substrate of thickness 1.6 mm and relative permittivity \( \varepsilon_r = 2.2 \), and are connected to the ground plane by metallic vias with diameter 0.7 mm. Also shown are the Cartesian and spherical reference coordinate systems, with origin at the local symmetry center. (a) Octagonal \( a_{\text{octagon}} = 7.6 \) mm, \( a_{\text{rhombus}} = 7.3 \) mm, total size: \( 70 \times 70 \) mm\(^2\). (b) Dodecagonal \( a_{\text{octagon}} = 7.6 \) mm, \( a_{\text{triangle}} = 7.1 \) mm, total size: \( 62 \times 62 \) mm\(^2\). (c) Penrose \( a_{\text{rhombus}} = 7.6 \) mm, \( a_{\text{rhombus}} = 7.2 \) mm, total size: \( 76 \times 76 \) mm\(^2\). (d) Periodic \( a_{\text{square}} = 7.6 \) mm, total size: \( 76 \times 76 \) mm\(^2\).

III. REPRESENTATIVE RESULTS

For finite-size HIS structures, the AMC behavior can be meaningfully ascertained by studying the radiation characteristics of a small electric dipole, laid parallel to the surface at a very close distance. Recalling basic image theory, the dipole will radiate very poorly when the surface is acting as standard (electric) ground plane, and much more efficiently when the surface is acting as an AMC, resulting in a low return-loss (RL) [2].

In our simulations, we placed a 16-mm-long electric dipole, parallel (\( z \)-directed) and very close to the surface \( \phi = y = 0, z = 0.7 \) mm, i.e., right above its local symmetry center, and computed via a full-wave commercial software package [20] the RL \( |S_{11}| \) spectra for the various configurations. Results are shown in Fig. 2 within the frequency range 4–12 GHz. For each geometry, one observes some (more or less pronounced) dips in the RL spectrum, corresponding to AMC-type behavior. In particular, the results for the octagonal HIS [Fig. 2(a)] agree fairly well with those observed in [18], with three main dips at frequencies \( 5, 6.5, \) and 10 GHz. Qualitatively similar behaviors, with less pronounced and frequency-shifted dips, are observed for the dodecagonal [Fig. 2(b)] and the Penrose [Fig. 2(c)] geometries, whereas the periodic HIS [Fig. 2(d)] turns out to exhibit more pronounced, but narrower, dips.

We then studied the radiation characteristics within the AMC bands. For conventional (metallic) ground planes, SWs can propagate along the metal-air interface and eventually radiate when reaching discontinuities (e.g., corners, edges), resulting in the appearance in the radiation pattern of ripples in the forward direction, and a degradation of the front-to-back ratio. It has been shown that periodically textured HIS ground-planes capable of strongly attenuating SW propagation can overcome these problems, thereby significantly improving the radiation pattern quality [2]. Fig. 3(a)–(d) shows some representative radiation patterns observed for the 16-mm-long electric dipole laid on the four HIS ground planes, at frequencies (marked as
arrows in Fig. 2) chosen so as to provide a reasonable tradeoff between high broadside directivity and low RL. As expected, one observes the lack of ripples in the forward direction and the low backward radiation ($\approx -15$ dB). Again, our results qualitatively confirm the rather high broadside directivity observed in [18] for the octagonal HIS [Fig. 3(a)], with half-power beamwidths of $52^\circ$ and $64^\circ$ in the principal planes $\phi = 0^\circ$ and $\phi = 90^\circ$, respectively. The lower directivity, as compared to that in [18], can be attributed to the smaller structure size considered here. Qualitatively similar results are observed for the dodecagonal [Fig. 3(b)] and Penrose [Fig. 3(c)] geometries. For the Penrose HIS, the chosen frequency (11 GHz) is considerably higher than for the other geometries; around lower-frequency dips, either the broadside directivity or the RL were found to be unsatisfactory. Moreover, the slight pattern asymmetry is attributable to the breaking of the underlying fivefold symmetry due to the presence of the dipole. For the periodic HIS in the considered parametric configuration [Fig. 3(d)], a moderately lower directivity was observed, with beamwidths of $90^\circ$ ($\phi = 0^\circ$ plane) and $70^\circ$ ($\phi = 90^\circ$ plane). Also, the RL is considerably higher than the aperiodic cases; at frequencies closer to the main RL dip [cf. Fig. 2(d)] the radiation patterns were found to exhibit a main-lobe splitting, with poorer broadside radiation (see, e.g., Fig. 4).

Finally, as in [18], we studied more in detail the SW propagation, for both transverse electric (TE) and magnetic (TM) polarizations. For the TE case, we placed two $y$-directed 16-mm-long electric dipoles laid 0.25 mm above the HIS, at two opposite ends, and computed the transmission ($\beta_{12}$) spectra. For the TM case, we considered two $z$-directed 8-mm-long monopoles. Results are shown in Fig. 5(a) and (b), respectively (with the simulation schematics illustrated in the insets), for the four geometries under study. In the TE case [Fig. 5(a)], the aperiodic HIS configurations exhibit several bandgaps characterized by strong SW attenuation, as observed in [18] for the octagonal geometry. Conversely, the periodic HIS exhibits only one bandgap with significant SW attenuation. In the TM case [Fig. 5(b)], one observes the presence of several dips and similar attenuation levels for all the geometries. These results were found to be rather robust with respect to variations in the probe positions.

IV. CONCLUSIONS AND OPEN ISSUES

In this letter, we have presented a comparative study of the EM properties of finite-size HIS ground-planes based on aperiodically ordered mushroom-type textures. We have considered three representative aperiodic-tiling geometries (octagonal, dodecagonal, and Penrose) and a standard periodic (square) geometry, and compared their performance as ground-planes for low-profile directive antennas, in terms of return-loss, radiation...
properties, and SW suppression. Our results, based on full-wave simulations, confirm those already observed in [18] (in connection with the octagonal geometry) and reveal the inherent suitability of more general aperiodic-tiling-based geometries to HIS applications. The availability of further and still largely unexplored degrees of freedom in the lattice geometry could open up new perspectives in the HIS design. Indeed, aperiodically ordered HIS structures seem potentially capable of yielding multiple AMC bands, strong attenuation of TE and TM SW propagation, and directive broadside radiation.

It should be emphasized that no effort was made at this stage to optimize the various parametric configurations, and to assess the actual superiority of aperiodically ordered HISs as compared to standard periodic configurations. In this connection, more comprehensive parametric studies and deeper understanding of the underlying phenomenologies are needed for a fruitful exploitation of the aforementioned geometry-related degrees of freedom in the engineering of bandgap and AMC properties. Among the open issues, currently under investigation, it is worth mentioning the following.

- The role of local order and symmetry. In this connection we are planning a comprehensive parametric study involving various randomly cut tiling realizations (not necessarily preserving the local symmetry center) of different sizes, and with different source positions and orientations. Also of interest, is the study of periodic Archimedean-type tiling geometries [21] featuring different (e.g., square- and equilateral-triangle-shaped) tiles, so as to separate the aperiodic effects from the presence of different tile shapes (see also the discussion in [18]).
- The study of possible mechanisms for performance control and optimization.
- The study of aperiodically ordered uniplanar HIS configurations.

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REFERENCES


